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Enantiopure Benzamidinate/Cyclooctatetraene Complexes of the Rare-Earth Elements: Synthesis, Structure, and Magnetism

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Supporting Information

ABSTRACT: Novel chiral rare-earth half sandwich complexes containing the chiral amidinate ligand, (S,S)-N,N'-bis(1-phenylethyl)benzamidinate ((S)-PEBA) or the 'Bu analogue (S,S)-N,N'-bis(1-phenylethyl) pivalamidinate ((S)-PETA) and the cyclooctatetraene dianion $(C_8H_8)^{2-}$ (COT) were synthesized. All complexes were fully characterized and the solid state structures were established by single-crystal crystallography. Magnetic property measurements indicated that the complex $[{(S)}-$ PETA}Er(COT)(THF)] is a typical field-induced single-molecule magnet (SMM), of which the magnetic properties are in reasonable agreement with high-level quantum chemical calculations. Due to predominantly electrostatic ligand field effects, the



relativistic J = 15/2 ground manifold of Er(III) is split into 8 Kramers doublets spanning a range of ~500 cm⁻¹, representable by means of higher order spin tensor operators. According to the quantum chemical calculations as well as from analysis of the experimental data, the magnetic relaxation pathways in these complexes are most likely to be of vibrational rather than electronic origin.

INTRODUCTION

Along with the widely established cyclopentadienyl ligand and its derivatives, the cyclooctatetraene dianion $(C_8H_8)^{2-}$ (COT) is one of the most important ligands in organo-f-element chemistry.¹⁻³ The large and flat cyclooctatetraene ligand represents a valuable alternative to the cyclopentadienyl ligands.^{4,5} Pioneering work on cyclooctatetraene lanthanide organometallic chemistry was reported by Streitwieser and coworkers in 1970.⁶ Since then, numerous lanthanide cyclooctatetraenide complexes with different structural motifs have been published.^{1–3} Among the longest known lanthanide COT complexes are symmetrical sandwich complexes containing $[Ln(C_8H_8)_2]^-$ anions.^{7–9} Also, triple-decker complexes with a COT middle deck $[(\mu - \eta^8 \cdot \eta^8 \cdot \text{COT}) \{ \text{Ln}(\eta^5 \cdot \text{Cp}^*) \}_2]^{10,11}$ (Cp* = η^5 -C₅Me₅) have been reported for a series of divalent lanthanides. Spectacular structures were obtained by using silyl-substituted COT ligands such as $[C_8H_6(SiMe_3)_2-1,4]^{2-}$ (= COT''),¹² $[C_8H_6(Si^{1}Pr_3)_2-1,4]^{2-}$ (= COT^{1,4SiiPr3}),¹³ and $[1,4-1,4]^{2-}$ $R_2C_8H_6]^{2-}$ (R = o-(dimethylsilyl)-N,N-dimethylaniline).^{8,14} These bulky COT ligands allowed for the synthesis of new double-decker and triple-decker sandwich-complexes, e.g., $[(\mu \eta^{8}:\eta^{8}-\text{COT}^{1,4\text{SiiPr3}}\{\text{LnCp}^{*}\}_{2}\}^{13}$ and $[(\text{COT}'')\text{Ln}(\mu-\eta^{8}:\eta^{8}-\text{COT}'')\text{Ln}(\text{COT}'')]$ (Ln = Ce, Nd, Sm).¹² Besides these examples, numerous half-sandwich complexes [Ln(COT)R] (R = monoanionic ligand) such as R = $\hat{C}l_{,0}^{6} I_{,1}^{15,16} O_{,3}SCF_{,1}^{15,16}$

OR, ¹⁷ SR, ¹⁸ BH₄, ¹⁹ (Ph₂P)₂N, ²⁰ aminotroponiminates, ⁵ 1,4diazadienes,²¹ and {CH(PPh₂NSiMe₃)₂,^{22,23} have been synthesized.

Lanthanide cyclooctatetraene complexes have also been used as catalysts for hydroamination reactions,^{22,23} but the most recent application is the use as single-ion magnets (SIMs).^{24,25} In 2011, Murugesu et al. reported the Dy sandwich complex, $[Dy(COT'')_{2}Li(THF)(DME)]$, which showed single-molecule magnet (SMM) behavior and unusual multiple relaxation modes.²⁶ Gao et al. synthesized the heteroleptic lanthanide complex $[ErCp^*(COT)]$, which displayed a butterfly-shaped hysteresis loop at 1.8 K persisting up to 5 K.²⁷ After that, several novel lanthanide COT and COT'' single ion complexes have been investigated in terms of their SIM behavior.²⁸⁻³² As pointed out by Gao et al., the local symmetry generated by the ligand-field plays a crucial role for the design of SIMs.^{27,33-3} The most spectacular advances in this area were reported recently by using rare-earth metal compounds, e.g., as shown last year when a Dy(III) metallocene exhibits a record high anisotropy barrier.^{39,40} Since sandwich complexes possessing interesting magnetic properties,⁴¹ we were interested to study the magnetic properties of a half-sandwich complex having a

Received: June 14, 2018 Published: October 17, 2018 second asymmetric ligand in the coordination sphere. Therefore, we synthesized a series of chiral lanthanide COT amidinate complexes. We employed two different chiral amidinate ligands, which were reported by us earlier in terms of catalytic applications.^{42–45} These ligands are the chiral amidinate ((*S*,*S*)-*N*,*N*'-bis(1-phenylethyl) benzamidinate (*S*)-PEBA; Scheme 1)^{42–44,46} and the ^tBu analogue (*S*,*S*)-*N*,*N*'bis(1-phenylethyl) pivalamidinate ((*S*)-PETA); Scheme 1).⁴⁵

Scheme 1. Chiral Amidines (S)-HPEBA and (S)-HPETA



DISCUSSION

Synthesis.^{47,48} The title compounds $[{(S)-PETA}Ln-(COT)(THF)]$ (Ln = Y (1a), Sm (1b), Dy (1c), Er (1d), and Yb (1e)) and $[{((S)-PEBA }Ln(COT)(THF)]$ (Ln = Y (2a), Dy (2b), Er (2c), and Yb (2d)), respectively, were synthesized by two-step reactions (Scheme 2). Anhydrous





LnCl₃ salts were first reacted with Li(S)-PETA or Li(S)-PEBA in 1:1 stoichiometric ratio at room temperature to give $[Ln((S)-PETA)Cl_2(THF)_n]^{45}$ or $[Ln((S)-PEBA)-Cl_2(THF)_n]^{49}$ in situ. Then, the resulting intermediates were mixed with freshly prepared K₂(COT) and stirred for 24 h at room temperature. Crude products 1a-1e and 2a-2d were crystallized from *n*-pentane or *n*-pentane/toluene (Scheme 2). All complexes show the expected colors: Sm complex 1b, red; Yb complexes 1e and 2d, green; Y complexes 1a and 2a, light yellow; Dy complexes 1c and 2b, yellow; Er complexes 1d and 2c, pink.

All compounds were characterized by IR spectroscopy and elemental analysis. NMR spectra were recorded only for diamagnetic compounds **1a** and **2a**. A strong peak at 1630–1590 cm⁻¹ assigned to the C–N vibration of the amidinate ligand, was observed in the IR spectra of all complexes. In the ¹H NMR spectra of **1a** and **2a**, a sharp singlet peak was observed at δ 6.71 (**1a**), 6.84 (**2a**) for the aromatic η^8 -COT ring. These peaks are slightly shifted in comparison to the typical singlet peak of the COT ring, e.g., in [{CH-

 $(PPh_2NSiMe_3)_2 Y(COT)]^{22} \delta$ 6.54 ppm. In the NMR spectrum a symmetric coordination of the (S)-PEBA ligand is seen. The signal of the Ar–CHMe proton of the (S)-PETA ligand in 1a is observed at δ 4.69 ppm, and the signal of the corresponding methyl group at δ 1.43 ppm. The Ar–CHMe proton of the (S)-PEBA ligand in 2a appears as a quartet at δ 3.68 ppm and the signal of the corresponding methyl groups is seen at δ 1.24 ppm. In comparison to the corresponding peaks in K(S)-PEBA,⁴⁶ all peaks are upfield shifted by about 0.5 ppm in the NMR spectrum. Further details of the complex structures were deduced from the 2D ¹H, ⁸⁹Y gHMBC NMR spectra. The ⁸⁹Y signals at 68.3 ppm (1a) and 73.6 ppm (2a) in the NMR spectrum show cross peaks to the OCH₂ group of THF molecule and the COT ring.

Single crystals suitable for X-ray diffraction of all the compounds except 1a were obtained by slowly evaporating a pentane solution. In contrast, single crystals of 1a were grown by layering pentane on a toluene solution of the complex. All complexes are mononuclear and have similar structures, although they crystallize in two different space groups.

Crystallographic data for 1a-2d are summarized in Tables S1 and S2. Compounds 1a,c,d crystallize in the acentric orthorhombic space group $P2_12_12_1$. The molecular structure of 1a is displayed as an example in Figure 1. In contrast,



Figure 1. Molecular structure of 1a in solid state, omitting the hydrogen atoms and solvent molecules. Selected bond lengths [Å] or angles [°]: (data for the compounds 1d,e are also given). 1a: Cg_(COT)-Y 1.8454(5), Y-N1 2.333(5), Y-N2 2.365(5), Y-O 2.421(5); $Cg_{(COT)}$ -Y-N1 139.7(2), $Cg_{(COT)}$ -Y-N2 138.9(2), N1-Y-N2 55.9(2), N1-C9-N2 111.1(5). 1d: Cg_(COT)-Er 1.8285(3), Er-N1 2.315(4), Er-N2 2.353(3), Er-O 2.399(3); Cg_(COT)-Er-N1 139.51(15), Cg_(COT)-Er-N2 138.80(15), N1-Er-N2 56.20(12), N1-C9-N2 110.9(3). 1e: $Cg_{(COT)}$ -Yb 1.8033(3), Yb-N1 2.323(5), Yb-N2 2.315(6), Yb-O 2.367(5); Cg_(COT)-Yb-N1 141.23(2), Cg_(COT)-Yb-N2 134.83(3), N1-Yb-N2 56.80(2), N1-C9-N2 110.0(5). Cg = ring centroid. Although, compounds 1b,c are crystallizing in another space group, the corresponding bonding parameters are given here for comparison: 1b: $Cg_{(COT)}1-$ Sm1 1.9487(8), Sm1-N1 2.425(5), Sm1-N2 2.447(5), Sm1-O1 2.484(4), Cg_(COT)2-Sm2 1.9337(4), Sm2-N3 2.411(5), Sm2-N4 2.443(5), Sm2-O2 2.506(4); Cg_(COT)1-Sm1-N1 144.4(2), Cg_(COT)1-Sm1-N2 134.18(15), N1-Sm1-N2 53.30(15), N1-C9-N2 108.8(5), N3-Sm2-N4 53.77(16), N3-C42-N4 109.9(5). 1c: Cg_(COT)-Dy1 1.8568(4), Dy1-N1 2.358(8), Dy1-N2 2.397(8), Dy1-O1 2.419(6), Cg_(COT)2-Dy2 1.8596(7), Dy2-N3 2.344(7), Dy2-N4 2.381(8), Dy2-O2 2.456(7); Cg_(COT)1-Dy1-N1 140.515(1), Cg_(COT)1-Dy1-N2 142.721(2), N1-Dy1-N2 55.4(3), N1-C9-N2 110.8(8), N3-Dy2-N4 55.2(3), N3-C42-N4 111.4(9).

compounds 1b and 1e crystallize in the monoclinic space group $P2_1$. For 1a, there is one molecule of the corresponding complex and one molecule of toluene in the asymmetric unit 1a (Figure 1). In contrast, there are two lanthanide molecules with one molecule of solvent in the asymmetric unit of 1b or 1c. For 1d and 1e, again only one molecule of the corresponding complex is seen in the asymmetric unit. In all complexes, the rare-earth ion is coordinated by two N atoms from the (S)-PETA ligand, one O atom from THF, and the eight carbon atoms of the aromatic n^{8} -COT ring. Thus, a three-legged piano stool conformation is formed. The bond distances and angles are in the expected range. Thus, the average Ln-C(_{COT}) length is 2.584(8) Å (1a), 2.673(6)Å (1b), 2.588(13) Å (1c), 2.571(6) Å (1d), and 2.565(8)Å (1e), while the distance of Ln to the centroid of the COT ring $(Cg_{(COT)})$ is 1.8454(5) Å (1a), 1.9407(5)Å (1b), 1.8582(5)Å (1c),1.8285(3)Å (1d), and 1.8033(3) Å (1e). Both values are comparable to those reported in the literature, e.g., for $[K(18-c-6)(THF)_2][Y(COT)_2].$

Due to the crystal quality of 2a, the collected diffraction data were not sufficient to solve the crystal structure. As it is difficult to get good quality crystal, for 2a only the unit cell was determined. Compound 2c is isostructural to (S)-PEBA complexes 2d and 2e. They crystallize in the acentric triclinic space group P1, each with two molecules of the complexes in the asymmetric unit. The differences in bond lengths and angles of all four compounds are a result of the lanthanide contraction.

As a representative example, the structure of **2c** is discussed in detail (Figure 2). In the solid-state structure, the expected η^8



Figure 2. Molecular structure of 2c in solid state, omitting the hydrogen atoms for clarity. Selected bond lengths [Å] or angles [°] (data for the isostructural compounds 2c and 2d are also given): 2c: Cg_(COT)1-Er1 1.8074(2), Er1-N1 2.338(5), Er1-N2 2.356(5), Er1-O1 2.376(4), Cg_(COT)2-Er2 1.7948(2), Er2-N3 2.336(5), Er2-N4 2.342(5), Er2-O2 2.380(4); Cg_(COT)1-Er1-N1 138.05(15), Cg_(COT)1-Er1-N2 139.69(15), N1-Er1-N2 57.23(2), N1-C9-N2 115.6(5), N3-Er2-N4 57.43(2), N3-C44-N4 114.4(5). 2d: Cg_(COT)1-Yb1 1.7861(3), Yb1-N1 2.338(7), Yb1-N2 2.322(7), Yb1-O1 2.360(6)Cg_(COT)2-Yb2 1.7690(3), Yb2-N3 2.319(7), Yb2-N4 2.310(7), Yb2-O2 2.364(6);Cg_(COT)1-Yb1-N1 139.1(2), Cg_(COT)1-Yb1-N2 137.5(2), N1-Yb1-N2 57.8(2), N1-C9-N2 115.5(7), N3-Yb2-N4 58.1(2), N3-C44-N4 115.8(7). Although compound 2b crystallizes in another space group, the corresponding bonding parameters are given here for comparison: Cg_(COT)-Dy 1.8443(3), Dy-N1 2.361(3), Dy-N2 2.401(3), Dy-O1 2.417(3); $Cg_{(COT)}$ -Dy-N1 137.01(10), $Cg_{(COT)}$ -Dy-N2 136.36(9), N1-Dy-N2 56.37(11), N1-C9-N2 115.4(3).

coordination of the COT ligand is seen. The planar η^8 -ring shows no significant distortion within the carbon framework (average C–C bond distances 1.4074 Å). The Er–Cg(_{COT})

distances are in the range between 2.551(7) and 2.606(7) Å, with an average value of 2.574 Å, which is in agreement with $[{CH(PPh_2NSiMe_3)_2}ErCOT)]$ (2.560(3)–2.608(3) Å, avg 2.579 Å).²² The (S)-PEBA group is almost symmetrically attached to the metal center (Er1–N1 2.338(5), Er1–N2 2.356(5), Er2–N3 2.336(5), and Er2–N4 2.342(5).

Dy complex **2b** crystallizes in the orthorhombic space group $P2_12_12_1$. There are four complex molecules in each unit cell. Its structure is similar to those of **2c** and **2d**. The Dy-Cg(_{COT}) distance (1.8443(3) Å), which is longer than that in [Dy(COT)(Cp*)] (1.6393(2) Å),⁵¹ is slightly shorter than that of [K(18-C-6)][Dy(COT)_2] (1.9092(3) Å).

Magnetic Properties. The magnetic properties of Dy and Er compounds 1c, 2b,c, and 1d were measured on polycrystalline samples. As shown in Figure 3, the χT products at 300 K



Figure 3. Plots of the χT vs T for complexes 1c, Dy, and 1d, Er, at 1000 Oe dc field.

of 1c and 1d are 13.88 and 11.00 cm³ K mol⁻¹, respectively, close to the expected value of 14.17 cm³ K mol⁻¹ for one uncoupled Dy(III) metal ion (J = 15/2, g = 4/3) and 11.48 $cm^3 K mol^{-1}$ for one uncoupled Er(III) metal ion (J = 15/2, g = 6/5), respectively. Upon cooling, the χT products of both 1c and 1d decrease slowly between 300 to 100 K, then decrease more rapidly to reach 10.46 and 7.85 cm³ K mol⁻¹ at 1.8 K, respectively. The continuous decrease of the χT product on lowering the temperature can be explained as result of the progressive depopulation of m_I sublevels and magnetic anisotropy of Dy(III) and Er(III) ions.^{50,52} The corresponding magnetic data for compounds 2b and 2c are shown in Figure S30. The field dependence of magnetization of 1c and 1d from zero dc field to 7 T at 2, 3, and 5 K are shown in Figure 4, and the corresponding maximum values reached are 5.42 and 4.38 $\mu_{\rm B}$. The lack of saturation of magnetization at 7 T can likely be attributed to the inherent magnetic anisotropy of the metal ions.^{53–56}

In order to verify their potential SMM behavior alternating current (ac) magnetic susceptibility studies were carried out on freshly filtered samples of 1c, 1d, 2b, and 2c. The magnetization relaxation of all systems was probed under zero dc field using ac susceptibility measurements as a function of the temperature at 1000 Hz. None of the complexes show an out-of-phase signal without the application of a dc field (Figure S25). One reason for this finding could be the very fast quantum tunneling of the magnetization (QTM) commonly



Figure 4. Plots of M vs H at 2, 3, and 5 K for complexes 1c, Dy, (top), and 1d, Er, (bottom). Lines are guides for the eye.

seen in pure lanthanide complexes.^{57,58} In order to suppress QTM, ac susceptibility measurements in the presence of weak dc fields were performed (500-3000 Oe) at 1.8 K (Figure S26). For complexes 1c, 2b, and 2c, no maximum in the out of phase (χ'') signal was observed even with an applied dc field indicating the absence of SMM behavior for these complexes, in contrast to 1d, which exhibits slow relaxation of magnetization under applied static fields. The most effective suppression of QTM was observed at H_{dc} = 1000 Oe. Therefore, ac susceptibility measurements for 1d were carried out at this field strength (Figures 5, 6, and S27) revealing the typical behavior for a field-induced SMM. Cole-Cole plots of χ' versus χ'' , between 2.8 and 7.2 K (Figure S28) revealed semicircular profiles and were fitted to a generalized Debye model with α ranging from 0.01 to 0.12. This indicates a narrow distribution of relaxation processes in complex 1d. To extract the associated relaxation parameters,⁵⁹ an Arrhenius plot of the relaxation time (τ) against 1/T was constructed (Figure 7). The high temperature region was fitted to a straight line giving an Orbach energy barrier of 49.7 K and a relaxation time of 1.8×10^{-7} s. Owing to the significant deviation from linear behavior at low temperatures, the following eq 1 was used to analyze the relaxation data over the whole temperature range:



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Figure 5. Plots χ'' vs temperature for complex 1d, Er, with a dc field of 1000 Oe.



Figure 6. Plots χ'' vs frequency for complex 1d, Er, with a dc field of 1000 Oe.



Figure 7. Plots of $\ln(\tau)$ vs T^{-1} for **1d** under 1000 Oe dc field. Circles represent experimental data. The different lines represent different fits that were used to analyze the data.

$$\tau^{-1} = AT + B + CT^{n} + \tau_{0}^{-1} e^{-E_{\rm eff}/k_{\rm B}T}$$
(1)

The first two terms (AT and B) describe the direct relaxation and QTM processes, the third term (CT^n) describes a Raman process, and the fourth term ($\tau_0^{-1} e^{-E_{\rm eff}/k_{\rm B}T}$) describes an Orbach process. The results of the fit are given in Table S12. Although the fit using eq 1 is slightly better, the complete Arrhenius diagram can be fitted using only the Raman term of eq 1 indicating that this is the dominant relaxation mechanism. Because of the constrained nature of metal $-\eta^8$ -COT vibrations,⁴⁰ we conjecture that a major relaxation pathway comprises the THF and (S)-PEBA ligands. Furthermore, it cannot be excluded that electronic effects, like the reduced electron density of the N-ligands resulting from substituting the ^tBu with a phenyl group or other ligand field effects could facilitate dissipation in **2b** and **2c**.

Calculated Magnetic Data. Quantum chemical calculations were performed to obtain zero-field splittings as well as macroscopic magnetic properties for Er(III) compounds 1d and 2c, as well as the Yb(III) 1e and Dy(III) derivatives 1c (see the "Quantum Chemical Calculations" section). The calculated magnetic susceptibilities and field-dependent magnetization show reasonable agreement with the experimental data (Figures 8 and S30, as well as Tables S4 and S5, respectively). The calculated values of the susceptibilities occur mainly in the low temperature range. Here, it has to be recognized that only isolated molecules are considered in the calculations.



Figure 8. Comparison of the calculated and measured susceptibilities (top) and the field dependence of the magnetization at 2, 3, and 5 K for **1d**.

Considering the field-induced Er(III) based SMM, 1d, the calculated magnetization is very similar to the experimental data for small fields and increases to slightly larger values for higher fields. The energy spectrum of the eight lowest Kramers doublets (Table 1) stemming from the ${}^{4}I_{15/2}$ ground state of

Table 1. Energies (in cm^{-1}) of the Lowest Electronic States of 1d without and with Inclusion of Spin–Orbit Coupling^{*a*}

	CASSCF/CI ^b	CASPT2 ^b		SOCI ^c	SOCI* ^c
1	0 (0)	221 (119)	1	0 (0)	0
2	9 (12)	0 (0)	2	64 (50)	124
3	42 (37)	511 (231)	3	110 (102)	186
4	81 (75)	482 (302)	4	131 (130)	212
5	110 (109)	248 (219)	5	185 (187)	283
6	143 (147)	130 (198)	6	260 (271)	331
7	228 (237)	395 (353)	7	374 (390)	439
8	290 (306)	362 (371)	8	435 (459)	590
9	332 (351)	542 (476)			
10	391 (413)	581 (500)			
11	393 (415)	490 (477)			
12	497 (534)	811 (680)			
13	498 (534)	801 (680)			

^{*a*}CAS(11,7)SCF/CI level and CAS(11,7)PT2 level; CAS(11,14) values in parentheses. ^{*b*}The CASSCF/CASPT2 states correspond to the ⁴I term of the free ion in presence of the ligand field). ^{*c*}Spin—orbit coupling (SOCI and SOCI*) within the full active space (only the first 32 microstates in the (11,14) case) results in eight low lying Kramers doublets which correspond to the J = 15/2 ground state of the ⁴I. The next states follow at 6600 cm⁻¹.

the free ion indicate a significant rhombicity induced by the asymmetric ligand field. This is in line with the experimental behavior, namely, the low value of the Orbach barrier compared to the energy of the first excited state and the fast relaxation observed for zero dc field. Owing to the asymmetry of the ligand field, the low energy spectrum is insufficiently described using the well-established second-order description involving the D-tensor. However, a reasonable agreement is achieved using higher (fourth to eighth) order (extended Stevens) operator equivalents (Table S6) comprising terms of lower symmetry, e.g., S_zS_+ . Considering all eight Kramers doublets (J = 15/2) an isotropic g-tensor with g = 1.19identical to the value of the free ion was obtained. For the first Kramers doublet and an effective spin of S = 1/2, an axial gtensor with $g_z = 17.2$ and g_x and g_y close to zero is observed. The values for the other low lying Kramers doublets are listed in Table S7

The electronic and thus magnetic properties of the other Er(III) based complex 2c are rather similar to those obtained for 1d (see Tables S7 and S8 and Figure S30). However, since 2c does not show slow magnetic relaxation, we conjecture that vibrational rather than electronic effects dominate magnetic relaxation in 1d.

The magnetic structures of the non-SMM complexes based on Yb(III) and Dy(III) do not display distinct axial behavior (Table S10). As in 1d, relaxation via electronically excited states is not the dominant channel for both 1c and 1e because of the magnitude of the first excitation energy (82 cm⁻¹ for 1c, 334 cm⁻¹ for 1e).

Thus, based on the results of our quantum chemical calculations, a fast relaxation via spin-phonon coupling is

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more likely than Orbach-type processes, although other effects facilitating quantum tunneling effects cannot be excluded.

CONCLUSION

A series of chiral lanthanide complexes based on (S)-PEBA/(S)-PETA and COT has been synthesized and the compounds fully characterized. A study of the magnetic properties of Dy and Er compounds 1c and 1d indicated that complex 1d is a typical field-induced SIM. Quantum chemical calculations on the electronic states corresponding to the ${}^{4}I_{15/2}$ ground term states reveal the presence of an axial g-tensor for the lowest Kramers doublet. For the quantitative description of the ligand field splitting, higher order Stevens operators are necessary in order to account for the inherent ligand field asymmetry. Due to the temperature dependence of the magnetic relaxation data, and supported by quantum chemical calculations, Raman-type relaxation mechanisms are suggested to play a dominant role in the asymmetric lanthanide complexes investigated.

EXPERIMENTAL SECTION

General Methods. All air- and water-sensitive materials were prepared under nitrogen atmosphere by using either Schlenk line or glovebox. THF was dried by distilling from potassium benzophenone ketyl under nitrogen before use. Toluene and n-pentane were dried with a MBraun solvent purification system (SPS-800). Deuterated solvents were purchased from Carl Roth GmbH (99.5 atom % D) and were dried and stored under vacuum with Na/K alloy. NMR spectra were recorded on a Bruker Avance II 300 MHz or Bruker Avance III 400 MHz. Chemical shifts are referenced to internal solvent resonances and are reported relative to tetramethylsilane (1H and ¹³C NMR) and [Y(NO₃)₃] (⁸⁹Y NMR), respectively. Elemental analyses were carried out with an Elementar Vario Micro Cube. IR spectra were performed on a Bruker TENSOR 37 spectrometer via the attenuated total reflection method (ATR). LnCl₃,⁶⁰ lithium-N,N'bis((S)-1-phenylethyl)benzamidinate (Li(S)-PEBA),^{46,61} lithium-N,N'-bis((S)-1-phenylethyl)*tert*-butylamidinat (Li(S)-PETA),⁶ and $K_2C_8H_8$ $(K_2COT)^{62}$ were prepared following the literature procedures. The other chemicals were commercially available and used without further purification.

Magnetic susceptibility data (1.8–300K) were collected on powdered samples using a Quantum Design model MPMS-XL SQUID instrument under a 1000 Oe applied magnetic field. Magnetization isotherms were collected at 2, 3 and 5 K between 0 and 7T. The ac susceptibility measurements were carried out under an oscillating ac field of 3 Oe and ac frequencies ranging from 1 to 1500 Hz. Data were corrected for diamagnetism using Pascal constants and a sample holder correction.

Synthesis of [{(S)-PETA}Ln(COT)(THF)] [Ln = Y(1a), Sm (1b), Dy (1c), Er (1d), Yb (1e)].⁴⁸ Anhydrous LnCl₃ (0.63 mmol) was treated with Li(S)-PETA (200 mg, 0.63 mmol) in 10 mL of THF under a nitrogen atmosphere. The reaction was allowed to warm up to room temperature and kept stirring for 1 h. Then, 10 mL of THF solution of freshly prepared K₂COT was added, and immediately, a color change was observed. After stirring for an additional 24 h. THF was removed in vacuo, and 20 mL of *n*-pentane was added. The solid was removed by filtration and filtrate was concentrated. The residue was purified by crystallization from either *n*-pentane or toluene/*n*pentane (1:2) (1b).

1a. Yellow crystals. Yield: 142 mg (based on single crystals), 0.248 mmol, 39%. ¹H NMR (400.13 MHz, C₆D₆, 298 K): δ (ppm) = 7.38 (d, *J* = 7.2 Hz, 4H, Ph), 7.25 (t, *J* = 7.6 Hz, 4H, Ph), 7.15–7.10 (m, 2H), 6.71 (s, 8H, COT), 4.69 (q, *J* = 5.8 Hz, 2H, CH), 3.24–3.14 (m, 2H, CH₂(THF)), 3.07–3.02 (m, 2H, CH₂(THF)), 1.42 (d, *J* = 5.8 Hz, 6H, CH₃), 0.95 (s, 9H, CH₃), 0.93–0.86 (m, 4H, CH₂(THF)). ¹³C{¹H} NMR (100.6 MHz, C₆D₆, 298 K): δ (ppm) = 178.4 (NCN), 148.9 (*i*-Ph), 128.2 (Ph), 126.3 (Ph), 126.1 (Ph), 94.2 (COT), 69.7 (CH₂(THF)), 55.4 (CH), 28.8 (CCH₃), 25.8 (NCHCH₃), 24.7

(CH₂(THF)), 23.0 (C(CH₃)₃). ⁸⁹Y{¹H} NMR (19.61 MHz, C₆D₆): δ (ppm) = 73.6. Anal. Calcd (%) for C₃₃H₄₃YN₂O (572.62): C, 69.22; H, 7.57; N, 4.89. Found: C, 69.86; H, 7.29; N, 4.42. IR (ATR) ν/cm^{-1} : 2953 (s), 2922 (s), 1653 (s), 1635 (s), 1490 (s), 1476 (s), 1447 (s), 1065 (m), 1026 (m), 897 (w), 757 (s), 698 cm⁻¹ (vs).

1b. Red crystals. Yield: 132 mg (based on single crystals), 0.208 mmol, 33%. Anal. Calcd (%) for $C_{33}H_{43}SmN_2O$ (634.07): C, 62.51; H, 6.84; N, 4.42. Found: C, 62.28; H, 6.71; N, 4.28. IR (ATR) $\nu/$ cm⁻¹: 2961 (s), 2923 (s), 1633 (s), 1490 (s), 1476 (s), 1448 (s), 1364 (m), 1253 (m), 1190 (m), 1066 (m), 1020 (m), 756 (s), 697 (vs), 542 cm⁻¹ (s).

1c. Yellow crystals. Yield: 171 mg (based on single crystals), 0.265 mmol, 42%. Anal. Calcd (%) for $C_{66}H_{86}Dy_2N_4O_2$ (1292.39): C, 61.34; H, 6.71; N, 4.34. Found: C, 61.07; H, 6.45; N, 4.20. IR (ATR) ν/cm^{-1} : 2961 (s), 2923 (s), 2856 (s), 1602 (s), 1476 (s), 1448 (s), 1397 (s), 1364 (m), 1278 (m), 1191 (m), 1066 (m), 1025 (m), 938 (w), 750 (s), 698 (vs), 636 (w), 510 cm^{-1} (s).

1d. Pink crystals. Yield: 111 mg (based on single crystals), 0.17 mmol, 27%. Anal. Calcd (%) for Anal. Calcd (%) for $C_{29}H_{35}ErN_2$ (578.87) (= 1d – THF): C, 60.17; H, 6.09; N, 4.84. Found: C, 60.09; H, 6.60; N, 4.43. IR (ATR) ν/cm^{-1} : 2963 (s), 1653 (s), 1635 (s), 1490 (s), 1447 (m), 1400 (m), 1362 (m), 1299 (m), 1186 (m), 1149 (m), 1065 (m), 1027 (s), 865 (m), 758 (s), 698 (vs), 672 cm⁻¹ (s).

1e. Green crystals. Yield: 96 mg (based on single crystals), 0.146 mmol, 23%. Anal. Calcd (%) for $C_{33}H_{43}YbN_2O$ (656.73): C, 60.35; H, 6.60; N, 4.27. Found: C, 60.01; H, 6.395; N, 4.14. IR (ATR) $\nu/$ cm⁻¹: 2963 (s), 2923 (s), 1653 (s), 1633 (s), 1559 (w), 1541 (w), 1490 (s), 1476 (s), 1448 (s), 1191 (m), 1066 (m), 1025 (m), 756 (s), 698 (vs), 543 cm⁻¹ (s).

Synthesis of [{(S)-PEBA}Ln(COT)(THF)] [Ln = Y(2a), Dy (2b), Er (2c), Yb(2d).⁴⁷ The synthesis of [{(S)-PEBA}Ln(COT)(THF)] was carried out following the similar method as that for [{(S)-PETA}Ln(COT)(THF)].

2a. Yellow crystals. Yield: 88 mg (based on single crystals), 0.148 mmol, 25%. ¹H NMR (400.13 MHz, C_6D_6 , 298 K): δ (ppm) = 7.14 (d, 6H, Ph), 7.11–7.01 (m, 4H, Ph), 6.92–6.90 (m, 3H, Ph), 6.83 (s, 8H, COT), 6.54–6.51 (m, 2H, Ph), 3.68 (m, J = 6.7 Hz, 1.7 Hz, 2H, CH), 3.10–3.03 (m, 4H, CH2 (THF)), 1.23 (d, J = 6.7 Hz, 6H, CH₃), 1.00–0.96 (m, 4H, CH₂(THF)). ¹³C{¹H} NMR (100.6 MHz, C_6D_6 , 298 K): δ (ppm) = 177.3 (NCN), 148.4 (*i*-Ph), 134.7 (*i*-Ph), 128.2 (Ph), 127.8 (Ph), 127.5 (Ph), 126.5 (Ph), 126.2 (Ph), 126.0 (Ph), 94.3 (COT), 69.8 (CH₂(THF)), 56.2 (CH), 25.1 (CH₃), 24.9 (CH₂(THF)). ⁸⁹Y{¹H} NMR (19.61 MHz, C_6D_6): δ (ppm) = 68.3. Anal. Calcd (%) for $C_{35}H_{39}YN_2O$ (591.58): C, 70.94; H, 6.63; N, 4.73. Found: C, 70.25; H, 6.64; N, 4.67. IR (ATR) ν/cm^{-1} : 3407 (m), 2955 (w), 2919 (w), 1637 (s), 1483 (s), 1451 (s), 1413 (s), 1309 (w), 1142 (m), 1072 (m), 1021 (m), 759 (s), 695 (vs), 541 cm⁻¹ (s).

2b. Yellow crystals. Yield: 86 mg (based on single crystals), 0.129 mmol, 22%. Anal. Calcd (%) for $C_{35}H_{39}DyN_2O$ (666.18): C, 63.10; H, 5.90; N, 4.20. Found: C, 62.67; H, 5.98; N, 4.00. IR (ATR) $\nu/$ cm⁻¹: 3300 (m), 2919 (w), 1700 (s), 1443 (s), 1373 (s), 1309 (w), 1298 (m), 1143 (m), 1072 (m), 1008 (m), 765 (s), 679 (vs), 477 cm⁻¹ (s).

2c. Pink crystals. Yield: 115 mg (based on single crystals), 0.171 mmol, 29%. Anal. Calcd (%) for $C_{35}H_{39}ErN_2O$ (670.94): C, 62.65; H, 5.86; N, 4.18. Found: C, 63.07; H, 6.09; N, 4.17. IR (ATR) ν/cm^{-1} : 3407 (w), 2955 (w), 2920 (w), 1637 (s), 1483 (s), 1446 (s), 1414 (s), 1365 (w), 1308 (w), 1268 (w), 1143 (w), 1072 (w), 1022 (m), 759 (s), 736(m), 695 (vs), 601 (w), 573 (w), 541 cm⁻¹ (s).

2d. Green crystals. Yield: 129 mg (based on single crystals), 0.191 mmol, 32%. Anal. Calcd (%) for $C_{31}H_{31}YbN_2$ (604.64) (= **2d** – THF): C, 61.58; H, 5.17; N, 4.63. Found: C, 61.00; H, 5.69; N, 3.92. IR (ATR) ν/cm^{-1} : 3407 (m), 2955 (w), 2919 (w), 1636 (vs), 1482 (s), 1451 (s), 1309 (w), 1298 (w), 1268 (w), 1142 (m), 1072 (m), 1028 (m), 765 (s), 698 (vs), 601 (s), 572 (s), 546 cm⁻¹ (s).

X-ray Crystallographic Studies of 1a–1e and 2a–2d. A suitable crystal was covered in mineral oil (Aldrich) and mounted on

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a glass fiber. The crystal was transferred directly to a cold stream of a STOE IPDS 2 or STOE StadiVari diffractometer.

All structures were solved by using the program SHELXS/T^{63,64} using Olex2.⁶⁵ The remaining non-hydrogen atoms were located from successive difference Fourier map calculations. The refinements were carried out by using full-matrix least-squares techniques on F^2 by using the program SHELXL.^{63,64} In each case, the locations of the largest peaks in the final difference Fourier map calculations, as well as the magnitude of the residual electron densities, were of no chemical significance. Positional parameters, hydrogen atom parameters, thermal parameters, and bond distances and angles have been deposited in the Supporting Information.

Data collection parameters are given in Table S1. Crystallographic data (excluding structure factors) for the structures reported in this paper have been deposited with the Cambridge Crystallographic Data Centre and the relevant codes are 1847599–1847607. Copies of the data can be obtained free of charge on application to CCDC, 12 Union Road, Cambridge CB21EZ, UK (fax: (+(44)1223–336–033; email: deposit@ccdc.cam.ac.uk).

Quantum Chemical Calculations.⁴⁷ For complexes 1c-e and 2c, quantum chemical calculations were performed to obtain the low energy states and the magnetic properties. All calculations are based on the experimental X-ray structure and were performed with the MOLCAS⁶⁶ and DIRAC⁶⁷ program packages if not mentioned elsewise. In the MOLCAS calculations, an ANO-RCC basis^{68–70} of triple- ζ quality was used for Dy, Yb, Er, N, and O, and a basis of double- ζ quality for C and H, the DIRAC calculations were performed on an Er embedded by a point charge model for the ligand field (see below) employing a dyall.v2z basis.⁷¹

Scalar relativistic effects were considered by the X2C operator.⁷² Because of the open shell character of the Ln(III) ions, wave-functionbased, multireference methods were applied: Energies and wave functions of the electronic states were therefore obtained by the nonrelativistic complete active space self-consistent field (CASSCF)⁷³ as well as the two-component relativistic Kramers-restricted complete active space configuration interaction method (KR-CASCI).⁷⁴

In the CASSCF calculations, the active space contained the seven 4f orbitals of Ln occupied with 9, 11, and 13 electrons. The orbitals were optimized for an averaged density constructed from all high spin states (21 sextets for CAS(9,11), 35 quartets for CAS(11,7), 7 doublets in CAS(13,7)) with equal weighting factors. The energies of the low spin states where obtained in a CASCI (CI: configuration interaction) calculation within the same active space. For 1d, dynamic correlation energies for the ground manifold were obtained by state-specific complete active space second-order perturbation theory (SS-CASPT2).⁷⁵ Spin orbit configuration interaction (SOCI) was included using quasi degenerate perturbation theory by coupling all electronic states in the active space for 1d, 1e, and 2c (147, 7, and 147 microstates). For 1c, an energy cutoff of 10 eV was used (535 microstates).

The calculations were performed with the RASSI module of MOLCAS using an atomic mean field approximation $(AMFI)^{76}$ for the spin–orbit integrals. The results are based on CASSCF/CI wave functions. In the SOCI calculations CASSCF/CI energies were taken, while in the SOCI* calculation for 1d, the SS-CASPT2 energies are used for the 13 lowest states. All other states were shifted by the average of the correlation energies. The magnetic properties such as magnetic susceptibility and field dependent magnetization, as well as g-tensors of the lowest eight Kramers doublets with a pseudospin of 1/2 and zero-field splitting parameters for a pseudospin of 7/2 (Yb(III)) and 15/2 (Er(III), Dy(III)) were obtained with the Single-ANISO module.^{77,78}

In order to check the influence of covalency of the ligands on the electronic structure of 1d, calculations on an electrostatic model were additionally performed where the electric field was constructed from point charges located at the positions of the atoms in the full complex. The values of the point charges were obtained from a Mulliken analysis⁷⁹ of the CAS(11,7)SCF ground states. The splitting of the low energy states is only slightly changed and the g-tensor of the lowest state remains axial (see the Supporting Information).

Concerning the complexity of the nonrelativistic electronic structure, in particular the encountered root-flipping in the CASPT2 calculations on 1d, the active space was enlarged by a second set of f-orbitals (CAS(11,14)SCF) in order to increase the flexibility of the reference wave functions and therefore improve the many-body perturbation theory treatment. The results are similar to the CAS(11,7) results confirming the root-flipping, but with a smaller separation of about 119 cm⁻¹. The effects on the electronic states are shown in Table 1 (values in parentheses). Because the effect of enlarging the active space on the electronic energies is considerably smaller than spin-orbit coupling, further calculations were dispensed. Instead, because of the importance of relativistic effects in lanthanide complex 1d, the validity of the SOCI approach was checked by means of two-component Kramers-restricted SCF and CI calculations using the DIRAC program suite (version 16).⁶⁷ In these calculations, spinorbit effects are considered in all basis functions as well as in the orbital optimization procedure, yielding a relativistic benchmark for the quasi-relativistic CASSCF/SOCI approach. The KR-SCF and KR-CI calculations employ a Kramers-restricted 2-spinor basis with an X2C operator for one-electron integrals together with AMFI corrections for the two-electron contributions and where performed on a point charge model (for a detailed description of the procedure, see Table S9). These calculations confirm the findings obtained from the quasirelativistic CAS/SOCI calculations on similar point charge field models performed with the SOCI program developed recently in Karlsruhe and Kaiserslautern (Table \$9).47 We therefore conclude that the chosen CAS(n_1 7)SCF/PT2 // SOCI method (n = 9, 11, 13) is adequate to describe the relativistic electronic and magnetic structure of the complexes under investigation.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.organo-met.8b00412.

Synthesis, IR, NMR, XRD, magnetic measurements, and quantum chemical calculations (PDF) Cartesian coordinates (XYZ)

Accession Codes

CCDC 1847599–1847607 contain the supplementary crystallographic data for this paper. These data can be obtained free of charge via www.ccdc.cam.ac.uk/data_request/cif, or by emailing data_request@ccdc.cam.ac.uk, or by contacting The Cambridge Crystallographic Data Centre, 12 Union Road, Cambridge CB2 1EZ, UK; fax: +44 1223 336033.

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Notes

The authors declare no competing financial interest.

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